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Positive and nodal single-layered solutions to supercritical elliptic problems above the higher critical exponents

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To Jean Mawhin on his 75th birthday, with great appreciation

ABSTRACT. We study the problem

 $-\Delta v + \lambda v = |v|^{p-2} v \text{ in } \Omega, \qquad v = 0 \text{ on } \partial\Omega,$

for $\lambda \in \mathbb{R}$ and supercritical exponents p, in domains of the form

 $\Omega := \{ (y, z) \in \mathbb{R}^{N-m-1} \times \mathbb{R}^{m+1} : (y, |z|) \in \Theta \},\$

where $m \geq 1$, $N - m \geq 3$, and Θ is a bounded domain in \mathbb{R}^{N-m} whose closure is contained in $\mathbb{R}^{N-m-1} \times (0, \infty)$. Under some symmetry assumptions on Θ , we show that this problem has infinitely many solutions for every λ in an interval which contains $[0,\infty)$ and p > 2up to some number which is larger than the $(m+1)^{st}$ critical exponent $2^*_{N,m} := \frac{2(N-m)}{N-m-2}$. We also exhibit domains with a shrinking hole, in which there are a positive and a nodal solution which concentrate on a sphere, developing a single layer that blows up at an m-dimensional sphere contained in the boundary of Ω , as the hole shrinks and $p \to 2^*_{N,m}$ from above. The limit profile of the positive solution, in the transversal direction to the sphere of concentration, is a rescaling of the standard bubble, whereas that of the nodal solution is a rescaling of a nonradial sign-changing solution to the problem

$$-\Delta u = |u|^{2_n^* - 2} u, \qquad u \in D^{1,2}(\mathbb{R}^n),$$

where $2_n^* := \frac{2n}{n-2}$ is the critical exponent in dimension n.

Keywords: Supercritical elliptic problem, positive solutions, nodal solutions, blow up, higher critical exponents.

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1. Introduction

We study the existence and concentration behavior of solutions to the problem

$$\begin{cases} -\Delta v + \lambda v = |v|^{p-2} v & \text{in } \Omega, \\ v = 0 & \text{on } \partial\Omega, \end{cases}$$
 (\$\mathcal{p}_p\$)

where Ω is a bounded domain in \mathbb{R}^N , $\lambda \in \mathbb{R}$, and p is supercritical, i.e., it is larger than the critical Sobolev exponent $2_N^* := \frac{2N}{N-2}$ for $N \geq 3$. We shall consider domains of the form

$$\Omega := \{ (y, z) \in \mathbb{R}^{N-m-1} \times \mathbb{R}^{m+1} : (y, |z|) \in \Theta \},$$

$$(1)$$

where $m \ge 1$, $N - m \ge 3$, and Θ is a bounded domain in \mathbb{R}^{N-m} whose closure is contained in $\mathbb{R}^{N-m-1} \times (0, \infty)$.

In domains of this type, the true critical exponent is $2_{N,m}^* := \frac{2(N-m)}{N-m-2}$, which is the critical Sobolev exponent in the dimension of Θ and is larger than 2_N^* . Indeed, one can easily verify that the solutions to the problem (\wp_p) which are radial in the variable z, correspond to the solutions of the problem

$$\begin{cases} -\operatorname{div}(f(x)u) + \lambda f(x)u = f(x) |u|^{p-2} u & \text{in } \Theta, \\ u = 0 & \text{on } \partial\Theta, \end{cases}$$
(2)

where $f(x_1, ..., x_{N-m}) = x_{N-m}^m$. Standard variational methods show that this last problem has infinitely many solutions for $p \in (2, 2_{N-m}^*)$, hence, also does the problem (\wp_p) . On the other hand, Passaseo showed in [18, 19] that, if $\lambda = 0$ and Θ is a ball centered on the half-line $\{0\} \times (0, \infty)$, then the problem (\wp_p) does not have a nontrivial solution for $p \geq 2_{N-m}^* = 2_{N,m}^*$. The number $2_{N,m}^*$ has been called the $(m+1)^{st}$ critical exponent in dimension N.

The concentration behavior of solutions to the problem (\wp_p) for $\lambda = 0$ and $p \in (2, 2^*_{N,m})$, as $p \to 2^*_{N,m}$ from below, has been investigated in several papers. In [10], del Pino, Musso and Pacard exhibited positive solutions which concentrate and blow up at a nondegenerate closed geodesic in $\partial\Omega$, as p approaches the second critical exponent $2^*_{N,1}$ from below. For any $m \ge 1$, positive and sign-changing solutions in domains of the form (1) were constructed in [1, 13]. These solutions concentrate and blow up at one or several m-dimensional spheres, as $p \to 2^*_{N,m}$ from below. In all of these cases the limit profile of the solutions, in the transversal direction to each sphere of concentration, is a sum of rescalings of $\pm U$, where

$$U(x) := [n(n-2)]^{(n-2)/4} \left(\frac{1}{1+|x|^2}\right)^{(n-2)/2}$$

is the standard bubble in dimension n := N - m, which is the only positive solution to the limit problem

$$-\Delta u = |u|^{2_n^* - 2} u, \qquad u \in D^{1,2}(\mathbb{R}^n), \tag{3}$$

up to translation and dilation.

It was recently shown in [4] that there exist nonradial sign-changing solutions to the problem (3), that do not resemble a sum of rescaled positive and negative standard bubbles, which occur as limit profiles for concentration of sign-changing solutions to the problem (\wp_p) that blow up at a single point, as $p \to 2_N^*$ from below. For the higher critical exponents $2_{N,m}^*$ with $m \ge 1$, it was shown in [5] that for every λ in some interval which contains $[0, \infty)$ there are sign-changing solutions to the problem (\wp_p) , in domains of the form (1), which concentrate and blow up at an *m*-dimensional sphere, as $p \to 2_{N,m}^*$ from below, whose limit profile in the transversal direction to the sphere of concentration is a nonradial sign-changing solution to (3), like those found in [4].

The study of concentration phenomena for p approaching 2_N^* from above, is a much more delicate issue, beginning with the fact that solutions to (\wp_p) for $p > 2_N^*$ do not always exist. For $\lambda = 0$, standard bubbles were used as basic cells in [8, 9, 16, 20] to construct positive solutions to the slightly supercritical problem (\wp_p) with $p = 2_N^* + \varepsilon$, for small enough $\varepsilon > 0$, in domains with a hole, using the Lyapunov-Schmidt reduction method. These solutions blow up, as $\varepsilon \to 0$, and their limit profile at each blow-up point is a rescaling of the standard bubble. Solutions in some contractible domains were constructed in [14, 15].

Quite recently, sign-changing solutions to the slightly supercritical problem (\wp_p) with $p = 2_N^* + \varepsilon$, $\varepsilon > 0$, were exhibited by Musso and Wei [17] in domains with a small fixed hole, and by Clapp and Pacella [6] in domains with a shrinking hole. The solutions obtained in [17] concentrate at two different points in the domain, as $\varepsilon \to 0$, and their limit profile at each of them is a rescaling of one of the sign-changing solutions to the limit problem (3) in \mathbb{R}^N constructed by del Pino, Musso, Pacard and Pistoia in [11], which resemble a large number of negative bubbles, placed evenly along a circle, surrounding a positive bubble, placed at its center. On the other hand, the sign-changing solutions exhibited in [6] concentrate at a single point in the interior of the shrinking hole, as the hole shrinks and $\varepsilon \to 0$, and their limit profile is a rescaling of a nonradial sign-changing solution to (3) like those found in [4].

For $m \geq 1$, the existence of solutions for $p = 2^*_{N,m} + \varepsilon$ and their concentration behavior seems to be, so far, an open question; see Problem 4 in [7]. In this paper we will show that, under some symmetry assumptions, the problem (\wp_p) has infinitely many solutions in domains of the form (1) for $p > 2^*_{N,m}$, up to some value which depends on the symmetries; see Theorem 2.3. We will also exhibit domains with a shrinking hole, in which there are positive and

sign-changing solutions which concentrate and blow up at an *m*-dimensional sphere contained in the boundary of Ω , as the hole shrinks and $p \to 2^*_{N,m}$ from above. The limit profile of the positive solutions, in the direction transversal to the sphere of concentration, will be a rescaling of the standard bubble, whereas that of the sign-changing ones will resemble one of the solutions to (3) that were found in [4].

We give, next, some examples of our results. For n := N - m, let B be an *n*-dimensional ball of radius δ_0 , centered on the half-line $\{0\} \times (0, \infty)$, whose closure is contained in the half-space $\mathbb{R}^{n-1} \times (0, \infty)$. We write the points in $\mathbb{R}^{n-1} \times (0, \infty)$ as (y, t) with $y \in \mathbb{R}^{n-1}$, $t \in (0, \infty)$ and we set

$$B_{\delta} := \{ (y,t) \in B : |y| > \delta \} \quad \text{if } \delta \in (0,\delta_0), \qquad B_0 := B,$$

$$\Omega_{\delta} := \{ (y,z) \in \mathbb{R}^{n-1} \times \mathbb{R}^{m+1} : (y,|z|) \in B_{\delta} \}, \qquad \Omega := \Omega_0.$$

We denote by O(k) the group of all linear isometries of \mathbb{R}^k and, for $v \in D^{1,2}(\mathbb{R}^N)$, we write

$$\|v\| := \left(\int_{\mathbb{R}^N} |\nabla v|^2\right)^{1/2}$$

The following results establish the existence of positive and sign-changing solutions to the problem (\wp_p) in Ω_{δ} and describe their limit profile as $\delta \to 0$ and $p \to 2^*_{N,m}$ from above. They are special cases of Theorems 2.3 and 4.4, which apply to more general domains, and are stated and proved in Sections 2 and 4, respectively.

THEOREM 1.1. There exists $\lambda_* \leq 0$ such that, for each $\lambda \in (\lambda_*, \infty) \cup \{0\}$, $\delta \in (0, \delta_0)$ and $p \in (2, \infty)$, the problem (\wp_p) has a positive solution $v_{\delta,p}$ in Ω_{δ} which satisfies

$$v_{\delta,p}(\gamma y, \varrho z) = v_{\delta,p}(y, z) \qquad \forall \gamma \in O(n-1), \ \varrho \in O(m+1), \ (y, z) \in \Omega_{\delta},$$

and has minimal energy among all nontrivial solutions to (\wp_p) in Ω_{δ} with these symmetries.

Moreover, there exist sequences (δ_k) in $(0, \delta_0)$, (p_k) in $(2^*_{N,m}, \infty)$, (ε_k) in $(0, \infty)$ and (ζ_k) in $B \cap [\{0\} \times (0, \infty)]$ such that

(i)
$$\delta_k \to 0$$
, $p_k \to 2^*_{N,m}$, $\varepsilon_k^{-1} \operatorname{dist}(\zeta_k, \partial \Theta) \to \infty$, and $\zeta_k \to \zeta$ with
 $\operatorname{dist}(\zeta, \mathbb{R}^{n-1} \times \{0\}) = \operatorname{dist}(B, \mathbb{R}^{n-1} \times \{0\}),$

(*ii*) $\lim_{k\to\infty} \left\| v_{\delta_k,p_k} - \widetilde{U}_{\varepsilon_k,\zeta_k} \right\| = 0$, where

$$\widetilde{U}_{\varepsilon_k,\zeta_k}(y,z) := \varepsilon_k^{(2-n)/2} U\left(\frac{(y,|z|)-\zeta_k}{\varepsilon_k}\right)$$

and U is the standard bubble in dimension n.

The number λ_* is negative if $m \geq 2$.

The solutions given by the previous theorem concentrate on an *m*-dimensional sphere, developing a positive layer which blows up at an *m*-dimensional sphere contained in the boundary of Ω and located at minimal distance to the plane of rotation $\mathbb{R}^{n-1} \times \{0\}$. The asymptotic profile of each layer in the transversal direction to its sphere of concentration is a rescaling of the standard bubble.

The next theorem gives sign-changing solutions to the problem (\wp_p) with a different type of asymptotic profile. For $n \geq 5$, we write $\mathbb{R}^{n-1} \equiv \mathbb{C}^2 \times \mathbb{R}^{n-5}$, and the points in \mathbb{R}^{n-1} as $y = (\eta, \xi)$, with $\eta = (\eta_1, \eta_2) \in \mathbb{C}^2$, $\xi \in \mathbb{R}^{n-5}$.

THEOREM 1.2. Assume that n = 5 or $n \ge 7$. Then, there exists $\lambda_* \le 0$ such that, for each $\lambda \in (\lambda_*, \infty) \cup \{0\}$, $\delta \in (0, \delta_0)$ and $p \in (2, 2^*_{N,m+1})$, the problem (\wp_p) has a nontrivial sign-changing solution $w_{\delta,p}$ in Ω_{δ} which satisfies

$$w_{\delta,p}(\eta,\xi,z) = w_{\delta,p}(\mathrm{e}^{\mathrm{i}\vartheta}\eta,\alpha\xi,\varrho z), \qquad w_{\delta,p}(\eta_1,\eta_2,\xi,z) = -w_{\delta,p}(-\bar{\eta}_2,\bar{\eta}_1,\xi,z),$$

for every $\vartheta \in [0, \pi)$, $\alpha \in O(n - 5)$, $\varrho \in O(m + 1)$ and $(y, z) \in \Omega_{\delta}$, and which has minimal energy among all nontrivial solutions to (\wp_p) in Ω_{δ} with these symmetry properties.

Moreover, there exist sequences (δ_k) in $(0, \delta_0)$, (p_k) in $(2^*_{N,m}, 2^*_{N,m+1})$, (ε_k) in $(0, \infty)$ and (ζ_k) in $B \cap [\{0\} \times (0, \infty)]$, and a nontrivial sign-changing solution W to the limit problem (3), such that

(i)
$$\delta_k \to 0$$
, $p_k \to 2^*_{N,m}$, $\varepsilon_k^{-1} \operatorname{dist}(\zeta_k, \partial \Theta) \to \infty$, and $\zeta_k \to \zeta$ with
 $\operatorname{dist}(\zeta, \mathbb{R}^{n-1} \times \{0\}) = \operatorname{dist}(B, \mathbb{R}^{n-1} \times \{0\}),$

- (ii) $W(\eta, \xi, t) = W(e^{i\vartheta}\eta, \alpha\xi, t)$ and $W(\eta_1, \eta_2, \xi, t) = -W(-\bar{\eta}_2, \bar{\eta}_1, \xi, t)$ for every $\vartheta \in [0, \pi)$, $\alpha \in O(n-5)$ and $(y, t) \in \mathbb{R}^{n-1} \times \mathbb{R} \equiv \mathbb{R}^n$, and W has minimal energy among all nontrivial solutions to (3) with these symmetry properties,
- (*iii*) $\lim_{k\to\infty} \left\| w_{\delta_k, p_k} \widetilde{W}_{\varepsilon_k, \zeta_k} \right\| = 0$, where

$$\widetilde{W}_{\varepsilon_k,\zeta_k}(y,z) := \varepsilon_k^{(2-n)/2} W\left(\frac{(y,|z|) - \zeta_k}{\varepsilon_k}\right)$$

The number λ_* is negative if $m \geq 2$.

The solutions given by the previous theorem concentrate on an *m*-dimensional sphere, developing a sign-changing layer which blows up at an *m*-dimensional sphere contained in the boundary of Ω and located at minimal distance to the plane of rotation $\mathbb{R}^{n-1} \times \{0\}$. The asymptotic profile of each layer in the transversal direction to its sphere of concentration is a rescaling of a nonradial sign-changing solution to the limit problem (3), like those found in [4].

As we mentioned before, the solutions to the anisotropic problem (2) give rise to solutions of the problem (\wp_p) in domains of the form (1). In Section 2 we will study a general anisotropic problem in an *n*-dimensional domain Θ . We will assume that Θ has some symmetries and we will establish the existence of infinitely many positive and sign-changing solutions to the anisotropic problem for supercritical exponents $p > 2_n^*$, up to some value which depends on the symmetries. These results extend those obtained in [6] for the problem with constant coefficients. In Section 3 we will describe the behavior of the minimizing sequences for the variational functional associated to the anisotropic problem for $p = 2_n^*$. These sequences, either converge to a solution, or they blow up. We will provide information on the location of the blow-up points and on the symmetries of the solutions to the limit problem (3) which occur as limit profiles. This will be used in Section 4 to obtain information on the concentration behavior and the limit profile of positive and sign-changing solutions to the problem (\wp_p) in domains with a shrinking hole, as the hole shrinks and $p \to 2^*_{N,m}$ from above.

2. Symmetries and existence for supercritical problems

Let Γ be a closed subgroup of O(n) and $\phi : \Gamma \to \mathbb{Z}_2$ be a continuous homomorphism of groups. A function $u : \mathbb{R}^n \to \mathbb{R}$ is said to be ϕ -equivariant if

$$u(\gamma x) = \phi(\gamma)u(x) \qquad \forall \gamma \in \Gamma, \ x \in \mathbb{R}^n.$$
(4)

If ϕ is the trivial homomorphism, then (4) simply says that u is a Γ -invariant function, whereas, if ϕ is surjective and u is not trivial, then (4) implies that u is sign-changing, nonradial and G-invariant, where $G := \ker \phi$.

Let Θ be a bounded Γ -invariant domain in \mathbb{R}^n , $n \geq 3$, and $a \in \mathcal{C}^1(\overline{\Theta})$, $b, c \in \mathcal{C}^0(\overline{\Theta})$ be Γ -invariant functions satisfying a, c > 0 on $\overline{\Theta}$. We assume that

there exists
$$x_0 \in \Theta$$
 such that $\{\gamma \in \Gamma : \gamma x_0 = x_0\} \subset \ker \phi$. (5)

This assumption guarantees that the space

$$D_0^{1,2}(\Theta)^{\phi} := \{ u \in D_0^{1,2}(\Theta) : u \text{ is } \phi \text{-equivariant} \}$$

is infinite dimensional; see [3]. As usual, $D_0^{1,2}(\Theta)$ denotes the closure of $\mathcal{C}_c^{\infty}(\Theta)$ in the Hilbert space

$$D^{1,2}(\mathbb{R}^n) := \{ u \in L^{2^*_n}(\mathbb{R}^n) : \nabla u \in L^2(\mathbb{R}^n, \mathbb{R}^n) \}$$

equiped with the norm

$$\|u\| := \left(\int_{\Theta} |\nabla u|^2\right)^{1/2}.$$

We shall also assume that the operator $-\operatorname{div}(a\nabla) + b$ is coercive in $D_0^{1,2}(\Theta)^{\phi}$, i.e., that

$$\inf_{\substack{u \in D_0^{1,2}(\Theta)^{\phi} \\ u \neq 0}} \frac{\int_{\Theta} (a(x) |\nabla u|^2 + b(x)u^2) \mathrm{d}x}{\int_{\Theta} |\nabla u|^2} > 0.$$
(6)

We set

$$||u||_{a,b}^{2} := \int_{\Theta} (a(x) |\nabla u|^{2} + b(x)u^{2}) \mathrm{d}x, \qquad |u|_{c;p}^{p} := \int_{\Theta} c(x) |u|^{p} \mathrm{d}x.$$

Assumption (6) implies that $\|\cdot\|_{a,b}$ is a norm in $D_0^{1,2}(\Theta)^{\phi}$ which is equivalent to $\|\cdot\|$. Note that, as c > 0, $|\cdot|_{c;p}$ is equivalent to the standard norm in $L^p(\Theta)$, which we denote by $|\cdot|_p$.

Our aim is to establish the existence of solutions to the problem

$$\begin{cases} -\operatorname{div}(a(x)\nabla u) + b(x)u = c(x)|u|^{p-2}u & \text{in }\Theta, \\ u = 0 & \text{on }\partial\Theta. \\ u(\gamma x) = \phi(\gamma)u(x), & \forall \gamma \in \Gamma, \ x \in \Theta, \end{cases}$$
(7)

for every 2 , where

$$d := \min\{\dim(\Gamma x) : x \in \overline{\Theta}\},\$$

 $\begin{array}{l} \Gamma x:=\{\gamma x:\gamma\in\Gamma\} \text{ is the }\Gamma\text{-orbit of }x,\,2_k^*:=\frac{2k}{k-2} \text{ if }k\geq 3 \text{ and }2_k^*:=\infty \text{ if }k=1,2. \text{ Note that }2_{n-d}^*>2_n^* \text{ if }d>0. \end{array}$

A (weak) solution to the problem (7) is a function $u \in D_0^{1,2}(\Theta)^{\phi} \cap L^p(\Theta)$ such that

$$\int_{\Theta} (a(x)\nabla u \cdot \nabla \psi + b(x)u\psi) dx - \int_{\Theta} c(x)|u|^{p-2}u\psi \, dx = 0 \qquad \forall \psi \in \mathcal{C}^{\infty}_{c}(\Theta).$$
(8)

Proposition 2.1 of [6] asserts that $D_0^{1,2}(\Theta)^{\phi}$ is continuously embedded in $L^p(\Theta)$ for any real number $p \in [1, 2^*_{n-d}]$, and that the embedding is compact for $p \in [1, 2^*_{n-d})$. The proof relies on a result by Hebey and Vaugon [12] which establishes these facts for Γ -invariant functions. Therefore, the functional

$$J_p(u) := \frac{1}{2} \left\| u \right\|_{a,b}^2 - \frac{1}{p} \left| u \right|_{c;p}^p$$

is well defined in the space $D_0^{1,2}(\Theta)^{\phi}$ if $p \in (2, 2^*_{n-d}]$.

LEMMA 2.1. For any real number $p \in (2, 2^*_{n-d}]$ the critical points of the functional J_p in the space $D_0^{1,2}(\Theta)^{\phi}$ are the solutions to the problem (7).

Proof. Let $u \in D_0^{1,2}(\Theta)^{\phi}$ be a critical point of J_p in $D_0^{1,2}(\Theta)^{\phi}$. Then,

$$J'_p(u)\vartheta = \int_{\Theta} (a(x)\nabla u \cdot \nabla \vartheta + b(x)u\vartheta - c(x)|u|^{p-2}u\vartheta) \,\mathrm{d}x = 0 \quad \forall \vartheta \in D^{1,2}_0(\Theta)^{\phi}.$$

As $D_0^{1,2}(\Theta)^{\phi} \subset L^p(\Theta)$ for $1 \leq p \leq 2^*_{n-d}$ we need only to prove that u satisfies (8). Let $\psi \in \mathcal{C}^{\infty}_c(\Theta)$, and define

$$\widetilde{\psi}(x) := \frac{1}{\mu(\Gamma)} \int_{\Gamma} \phi(\gamma) \psi(\gamma x) \mathrm{d}\mu,$$

where μ is the Haar measure on Γ . Note that $\tilde{\psi} \in D_0^{1,2}(\Theta)^{\phi}$. Observe also that, as u is ϕ -equivariant, we have that

$$\phi(\gamma)\nabla u(x) = \nabla \left(u \circ \gamma\right)(x) = \gamma^{-1}\nabla u(\gamma x) \qquad \forall \gamma \in \Gamma, \ x \in \Theta.$$

Since $J'_p(u)\widetilde{\psi}=0$, and a,b,c are Γ -invariant, using Fubini's theorem and performing a change of variable, we get

$$\begin{split} 0 &= \int_{\Theta} (a(x)\nabla u(x) \cdot \nabla \widetilde{\psi}(x) + b(x)u(x)\widetilde{\psi}(x) - c(x)|u(x)|^{p-2}u(x)\widetilde{\psi}(x))dx \\ &= \frac{1}{\mu(\Gamma)} \int_{\Theta} \int_{\Gamma} \left[a(x)\phi(\gamma)\nabla u(x) \cdot \gamma^{-1}\nabla\psi(\gamma x) + b(x)\phi(\gamma)u(x)\psi(\gamma x) \right. \\ &\quad -c(x)|\phi(\gamma)u(x)|^{p-2}\phi(\gamma)u(x)\psi(\gamma x)] d\mu dx \\ &= \frac{1}{\mu(\Gamma)} \int_{\Gamma} \int_{\Theta} \left[a(x)\gamma^{-1}\nabla u(\gamma x) \cdot \gamma^{-1}\nabla\psi(\gamma x) + b(x)u(\gamma x)\psi(\gamma x) \right. \\ &\quad -c(x)|u(\gamma x)|^{p-2}u(\gamma x)\psi(\gamma x)] dx d\mu \\ &= \frac{1}{\mu(\Gamma)} \int_{\Gamma} \int_{\Theta} \left[a(\gamma x)\nabla u(\gamma x) \cdot \nabla\psi(\gamma x) + b(\gamma x)u(\gamma x)\psi(\gamma x) \right. \\ &\quad -c(\gamma x)|u(\gamma x)|^{p-2}u(\gamma x)\psi(\gamma x)] dx d\mu \\ &= \frac{1}{\mu(\Gamma)} \int_{\Gamma} d\mu \int_{\Theta} \left[a(\xi)\nabla u(\xi) \cdot \nabla\psi(\xi) + b(\xi)u(\xi)\psi(\xi) \right. \\ &\quad -c(\xi)|u(x)|^{p-2}u(\xi)\psi(\xi)] d\xi \\ &= \int_{\Theta} \left[a(\xi)\nabla u(\xi) \cdot \nabla\psi(\xi) + b(\xi)u(\xi)\psi(\xi) - c(\xi)|u(x)|^{p-2}u(\xi)\psi(\xi) \right] d\xi. \end{split}$$

Therefore u is a solution to the problem (7).

The nontrivial critical points of the functional $J_p: D_0^{1,2}(\Theta)^{\phi} \to \mathbb{R}$ lie on the Nehari manifold

$$\mathcal{N}_{p}^{\phi} := \left\{ u \in D_{0}^{1,2}(\Theta)^{\phi} : \left\| u \right\|_{a,b}^{2} = \left| u \right|_{c;p}^{p}, \, u \neq 0 \right\},\$$

which is a \mathcal{C}^2 -Hilbert manifold, radially diffeomorphic to the unit sphere in $D_0^{1,2}(\Theta)^{\phi}$, and a natural constraint for this functional. Set

$$\ell_p^\phi := \inf\{J_p(u) : u \in \mathcal{N}_p^\phi\}.$$

Then, $\ell_p^{\phi} > 0$. A *least energy solution* to the problem (7) is a minimizer for J_p on \mathcal{N}_p^{ϕ} . The following result extends Theorem 2.3 in [6].

THEOREM 2.2. If $p \in (2, 2^*_{n-d})$ then the problem (7) has a least energy solution, and an unbounded sequence of solutions.

Proof. By Lemma 2.1, the critical points of the functional J_p in the space $D_0^{1,2}(\Theta)^{\phi}$ are the solutions to the problem (7). Proposition 2.1 of [6] asserts that $D_0^{1,2}(\Theta)^{\phi}$ is compactly embedded in $L^p(\Theta)$ for $p \in (2, 2^*_{n-d})$, hence, a standard argument shows that the functional $J_p: D_0^{1,2}(\Theta)^{\phi} \to \mathbb{R}$ satisfies the Palais-Smale condition. Therefore, J_p attains its minimum on \mathcal{N}_p^{ϕ} . Moreover, as the functional is even and has the mountain pass geometry, the symmetric mountain pass theorem [2] yields the existence of an unbounded sequence of critical values for J_p in $D_0^{1,2}(\Theta)^{\phi}$.

We now derive a multiplicity result for the supercritical problem (\wp_p) . Assume that the closure of Θ is contained in $\mathbb{R}^{n-1} \times (0, \infty)$ and, for $m \geq 1$, let

$$\lambda_1^{\phi} := \inf_{\substack{u \in D_0^{1,2}(\Theta)^{\phi} \\ u \neq 0}} \frac{\int_{\Theta} x_n^m |\nabla u|^2}{\int_{\Theta} x_n^m u^2}.$$
(9)

As the *n*-th coordinate x_n of x is positive for every $x \in \overline{\Theta}$, from the Poincaré inequality we obtain that $\lambda_1^{\phi} > 0$.

THEOREM 2.3. If $\lambda \in (-\lambda_1^{\phi}, \infty)$ and $p \in (2, 2^*_{n-d})$, then the problem (\wp_p) has a least energy solution and an unbounded sequence of solutions in

$$\Omega := \{ (y, z) \in \mathbb{R}^{n-1} \times \mathbb{R}^{m+1} : (y, |z|) \in \Theta \},\$$

which satisfy

$$v(\gamma y, \varrho z) = \phi(\gamma)v(y, z) \qquad \forall \gamma \in \Gamma, \ \varrho \in O(m+1), \ (y, z) \in \Omega.$$
(10)

Proof. A straightforward computation shows that v is a solution to the problem (\wp_p) in Ω which satisfies (10) if and only if the function u given by v(y, z) = u(y, |z|) is a solution to the problem (7) with $a(x) := x_n^m =: c(x)$ and $b(x) := \lambda x_n^m$. Moreover, v has minimal energy if and only if u does. Note that (6) is satisfied if $\lambda \in (-\lambda_1^{\phi}, \infty)$. So this result follows from Theorem 2.2.

For $p \in (2, 2^*_{n-d})$ let u_p be a least energy solution to the problem (7). Fix $q \in (2, 2^*_{n-d})$ and let $t_{q,p} \in (0, \infty)$ be such that $\tilde{u}_p := t_{q,p} u_p \in \mathcal{N}_q^{\phi}$, i.e.,

$$t_{q,p} = \left(\frac{\|u_p\|_{a,b}^2}{\|u_p\|_{c;q}^q}\right)^{\frac{1}{q-2}} = \left(\frac{\|u_p\|_{c;p}^p}{\|u_p\|_{c;q}^q}\right)^{\frac{1}{q-2}}.$$
(11)

We will show that $\lim_{p\to q} J_q(\tilde{u}_p) = \ell_q^{\phi}$. The proof is similar to that of Proposition 2.5 in [6]. We give the details for the reader's convenience, starting with the following lemma.

LEMMA 2.4. If $p_k, q \in (2, 2^*_{n-d}), p_k \to q$, and (u_k) is a bounded sequence in $D_0^{1,2}(\Theta)^{\phi}$, then

$$\lim_{k \to \infty} \int_{\Theta} \left(c(x) \left| u_k \right|^{p_k} - c(x) \left| u_k \right|^q \right) \mathrm{d}x = 0.$$

Proof. By the mean value theorem, for each $x \in \Theta$, there exists $q_k(x)$ between p_k and q such that

$$|u_k(x)|^{p_k} - |u_k(x)|^q| = |\ln |u_k(x)|| |u_k(x)|^{q_k(x)} |p_k - q|.$$

Fix r > 0 such that $[q - r, q + r] \subset (2, 2^*_{n-d})$. Then, for some positive constant C and k large enough,

$$|\ln |u_k|| |u_k|^{q_k} \le \begin{cases} \ln |u_k| |u_k|^{q+r} & \le C |u_k|^{2^*_{n-d}} & \text{if } |u_k| \ge 1, \\ \left(\ln \frac{1}{|u_k|}\right) |u_k|^{q-r} & \le C |u_k|^2 & \text{if } |u_k| \le 1. \end{cases}$$

As $D_0^{1,2}(\Theta)^{\phi}$ is continuously embedded in $L^p(\Theta)$ for $p \in [2, 2^*_{n-d}]$, we obtain

$$\left| \int_{\Theta} c\left(|u_{k}|^{p_{k}} - |u_{k}|^{q} \right) \right| \leq |c|_{\infty} \left(\int_{|u_{k}| \leq 1} ||u_{k}|^{p_{k}} - |u_{k}|^{q} | + \int_{|u_{k}| > 1} ||u_{k}|^{p_{k}} - |u_{k}|^{q} | \right)$$
$$\leq |c|_{\infty} C |p_{k} - q| \int_{\Theta} \left(|u_{k}|^{2} + |u_{k}|^{2^{*}_{n-d}} \right)$$
$$\leq \bar{C} |p_{k} - q| ||u_{k}||^{2^{*}_{n-d}}$$

for some positive constant \overline{C} , where $|c|_{\infty} := \sup_{x \in \Theta} |c(x)|$. Since (u_k) is bounded in $D_0^{1,2}(\Theta)$, our claim follows. PROPOSITION 2.5. For $q \in (2, 2^*_{n-d})$ we have that

$$\lim_{p \to q} \ell_p^{\phi} = \ell_q^{\phi}, \qquad \lim_{p \to q} t_{q,p} = 1, \qquad \lim_{p \to q} J_q\left(\widetilde{u}_p\right) = \ell_q^{\phi}.$$

Proof. Set

$$S_p^{\phi} := \inf_{u \in D_0^{1,2}(\Omega)^{\phi} \setminus \{0\}} \frac{\|u\|_{a,b}^2}{|u|_{c;p}^2}$$

It is easy to see that $\ell_p^{\phi} = \frac{p-2}{2p} \left(S_p^{\phi}\right)^{\frac{p}{p-2}}$. So, to prove the first identity, it suffices to show that $\lim_{p\to q} S_p^{\phi} = S_q^{\phi}$. From Hölder's inequality we get that $|u|_{c;q} \leq |c|_1^{(p-q)/pq} |u|_{c;p}$ if p > q. Hence, $S_q^{\phi} \geq |c|_1^{2(q-p)/pq} S_p^{\phi}$ if p > q. So, as p approaches q from the right, we have that

$$\limsup_{p \to q^+} S_p^\phi \le S_q^\phi.$$

Assume that $\liminf_{p \to q^+} S_p^{\phi} < S_q^{\phi}$. Then, there exist $\varepsilon > 0$ and sequences (p_k) in $(q, 2_{n-d}^*)$ and (u_k) in $D_0^{1,2}(\Omega)^{\phi}$ with $|u_k|_{c;p_k} = 1$ such that $||u_k||_{a,b}^2 < S_q^{\phi} - \varepsilon$. Lemma 2.4 implies that $\frac{||u_k||_{a,b}^2}{||u_k|_{c;q}^2} < S_q^{\phi}$ for k large enough, contradicting the definition of S_q^{ϕ} . This proves that

$$\lim_{p \to q^+} S_p^\phi = S_q^\phi$$

The corresponding statement when p approaches q from the left is proved in a similar way. Since $J_p(u_p) = \frac{p-2}{2p} \|u_p\|_{a,b}^2 = \ell_p^{\phi}$ we have that (u_p) is bounded in $D_0^{1,2}(\Omega)^{\phi}$ for p close to q. Lemma 2.4 applied to (11) yields $\lim_{p\to q} t_{q,p} = 1$. It follows that $\lim_{p\to q} J_q(\widetilde{u}_p) = \lim_{p\to q} \frac{q-2}{2q} \|t_{q,p}u_p\|_{a,b}^2 = \ell_q^{\phi}$, as claimed. \Box

3. Minimizing sequences for the critical problem

In this section we analize the behavior of the minimizing sequences for the problem (7) when p is the critical exponent $2_n^* = \frac{2n}{n-2}$. The solutions to the limit problem (3) will play a crucial role in this analysis. We denote the energy functional associated to (3) by

$$J_{\infty}(u) := \frac{1}{2} \|u\|^2 - \frac{1}{2^*} \|u\|_{2^*}^2$$

and, for any closed subgroup K of Γ , we set

$$D^{1,2}(\mathbb{R}^n)^{\phi|K} := \{ u \in D^{1,2}(\mathbb{R}^n) : u(\gamma z) = \phi(\gamma)u(z) \ \forall \gamma \in K, \ z \in \mathbb{R}^n \},$$
$$\mathcal{N}_{\infty}^{\phi|K} := \{ u \in D^{1,2}(\mathbb{R}^n)^{\phi|K} : u \neq 0, \ \|u\|^2 = |u|_{2^*}^{2^*} \},$$
$$\ell_{\infty}^{\phi|K} := \inf_{u \in \mathcal{N}_{\infty}^{\phi|K}} J_{\infty}(u).$$

If $K = \Gamma$ we write $\mathcal{N}^{\phi}_{\infty}$ and ℓ^{ϕ}_{∞} instead of $\mathcal{N}^{\phi|K}_{\infty}$ and $\ell^{\phi|K}_{\infty}$.

Recall that the Γ -orbit of a point $x \in \mathbb{R}^n$ is the set $\Gamma x := \{\gamma x : \gamma \in \Gamma\}$, and its isotropy group is $\Gamma_x := \{\gamma \in \Gamma : \gamma x = x\}$. Then, Γx is Γ -homeomorphic to the homogeneous space Γ/Γ_x . In particular, the cardinality of Γx is the index of Γ_x in Γ , which is usually denoted by $|\Gamma/\Gamma_x|$. If $\Gamma x = \{x\}$ then x is said to be a fixed point of Γ . We denote

$$\Theta^{\Gamma} := \{ x \in \Theta : x \text{ is a fixed point of } \Gamma \}.$$

For simplicity, we will write J_* , \mathcal{N}^{ϕ}_* and ℓ^{ϕ}_* instead of $J_{2^*_n}$, $\mathcal{N}^{\phi}_{2^*_n}$ and $\ell^{\phi}_{2^*_n}$.

THEOREM 3.1. Let (u_k) be a sequence in \mathcal{N}^{ϕ}_* such that $J_*(u_k) \to \ell^{\phi}_*$. Then, after passing to a subsequence, one of the following two possibilities occurs:

- 1. (u_k) converges strongly in $D_0^{1,2}(\Theta)$ to a minimizer of J_* on \mathcal{N}_*^{ϕ} .
- 2. There exist a closed subgroup K of finite index in Γ , a sequence (ζ_k) in Θ , a sequence (ε_k) in $(0, \infty)$ and a nontrivial solution ω to the problem (3) with the following properties:
 - (a) $\Gamma_{\zeta_k} = K$ for all $k \in \mathbb{N}$, and $\zeta_k \to \zeta$,
 - (b) $\varepsilon_k^{-1} \operatorname{dist}(\zeta_k, \partial \Theta) \to \infty \text{ and } \varepsilon_k^{-1} |\alpha \zeta_k \beta \zeta_k| \to \infty \text{ for all } \alpha, \beta \in \Gamma \text{ with } \alpha^{-1} \beta \notin K,$

(c)
$$\omega(\gamma z) = \phi(\gamma)\omega(z)$$
 for all $\gamma \in K$, $z \in \mathbb{R}^n$, and $J_{\infty}(\omega) = \ell_{\infty}^{\phi|K}$,

$$(d) \lim_{k \to \infty} \left\| u_k - \sum_{[\gamma] \in \Gamma/K} \phi(\gamma) \left(\frac{a(\zeta)}{c(\zeta)} \right)^{\frac{n-2}{4}} \varepsilon_k^{\frac{2-n}{2}} (\omega \circ \gamma^{-1}) (\frac{\cdot -\gamma\zeta_k}{\varepsilon_k}) \right\| = 0,$$

$$(e) \ \ell_*^{\phi} = \min_{x \in \overline{\Theta}} \frac{a(x)^{n/2}}{c(x)^{(n-2)/2}} \left| \Gamma/\Gamma_x \right| \ell_{\infty}^{\phi|\Gamma_x} = \frac{a(\zeta)^{n/2}}{c(\zeta)^{(n-2)/2}} \left| \Gamma/K \right| J_{\infty}(\omega).$$

Proof. The proof is exactly the same as that of Theorem 2.5 in [5], omitting

Let us state an interesting special case of Theorem 3.1.

the first two lines.

COROLLARY 3.2. Assume that every Γ -orbit in Θ is either infinite or a fixed point. Let (u_k) be a sequence in \mathcal{N}^{ϕ}_* such that $J_*(u_k) \to \ell^{\phi}_*$. Then, after passing to a subsequence, one of the following statements holds true:

1. (u_k) converges strongly in $D_0^{1,2}(\Theta)$ to a minimizer of J_* on \mathcal{N}_*^{ϕ} .

2. There exist a sequence (ζ_k) in Θ^{Γ} , a sequence (ε_k) in $(0, \infty)$ and a nontrivial ϕ -equivariant solution ω to the limit problem (3) such that $\zeta_k \to \zeta \in \overline{\Theta}$, $\varepsilon_k^{-1} \operatorname{dist}(\zeta_k, \partial \Theta) \to \infty, \ J_{\infty}(\omega) = \ell_{\infty}^{\phi}$,

$$\lim_{k \to \infty} \left\| u_k - \left(\frac{a(\zeta)}{c(\zeta)} \right)^{\frac{n-2}{4}} \varepsilon_k^{\frac{2-n}{2}} \omega \left(\frac{\cdot - \zeta_k}{\varepsilon_k} \right) \right\| = 0,$$

and

$$\frac{a(\zeta)^{n/2}}{c(\zeta)^{(n-2)/2}} = \min_{x \in \overline{\Theta^{\Gamma}}} \frac{a(x)^{n/2}}{c(x)^{(n-2)/2}}.$$

In particular, if every Γ -orbit in Θ has positive dimension, then (1) must hold true.

Proof. Since the group $K = \Gamma_{\zeta_k}$, given by case 2 of Theorem 3.1, has finite index in Γ and this index is the cardinality of the Γ -orbit of ζ_k , our assumption implies that $K = \Gamma$ and ζ_k is a fixed point. So case 2 of Theorem 3.1 reduces to case 2 of this corollary.

Note that the functions a and c determine the location of the concentration point ζ .

It was shown in [4, Theorem 2.3] that, if a = c = 1, b = 0 and $\Theta^{\Gamma} \neq \emptyset$, then ℓ_*^{ϕ} is not attained by J_* on \mathcal{N}_*^{ϕ} . So, if every Γ -orbit in $\Theta \setminus \Theta^{\Gamma}$ has positive dimension, statement 2 of Corollary 3.2 must hold true.

In the following section we will state a nonexistence result which allows us to obtain information on the limit profile of solutions to the problem (\wp_p) .

4. Blow-up at the higher critical exponents

Throughout this section we will assume that Θ is a Γ -invariant bounded smooth domain in \mathbb{R}^n whose closure is contained in $\mathbb{R}^{n-1} \times (0, \infty)$. Then, the points in $\{0\} \times (0, \infty)$ must be fixed points of Γ , so $\mathbb{R}^{n-1} \times \{0\}$ is Γ -invariant and we may regard Γ as a subgroup of O(n-1). We will also assume that $\Theta \setminus \Theta^{\Gamma}$ and Θ^{Γ} are nonempty, and that every Γ -orbit in $\Theta \setminus \Theta^{\Gamma}$ has positive dimension. As before, $\phi : \Gamma \to \mathbb{Z}_2$ will be a continuous homomorphism which satisfies assumption (5).

We set

$$\Theta_{\delta} := \{ x \in \Theta : \operatorname{dist}(x, \Theta^{\Gamma}) > \delta \} \text{ if } \delta > 0, \quad \text{and} \quad \Theta_0 := \Theta,$$

and we fix $\delta_0 > 0$ such that $\Theta_{\delta_0} \neq \emptyset$. For $m \ge 1$ and $\delta \in [0, \delta_0)$, we consider the problem

$$(\varphi_{\delta,p}^{\#}) \qquad \begin{cases} -\operatorname{div}(x_n^m \nabla u) + \lambda x_n^m u = x_n^m |u|^{p-2} u & \text{in } \Theta_{\delta}, \\ u = 0 & \text{on } \partial \Theta_{\delta}. \\ u(\gamma x) = \phi(\gamma) u(x), & \forall \gamma \in \Gamma, \ x \in \Theta_{\delta}, \end{cases}$$

where x_n^m denotes the function $x = (x_1, ..., x_n) \mapsto x_n^m$, and $\lambda \in (-\lambda_1^{\phi}, \infty)$, with λ_1^{ϕ} as defined in (9). Then, the operator $-\operatorname{div}(x_n^m \nabla) + \lambda x_n^m$ is coercive in $D_0^{1,2}(\Theta)^{\phi}$. So the data of this problem satisfy all assumptions stated at the beginning of Section 2.

Theorem 2.2 asserts that the problem $(\wp_{\delta,p}^{\#})$ has a least energy solution $u_{\delta,p}$ if $\delta \in (0, \delta_0)$ and $p \in (2, 2^*_{n-2})$, where

$$\mathfrak{d} := \min\{\dim(\Gamma x) : x \in \Theta \setminus \Theta^{\Gamma}\}.$$

Note that, by assumption, $\mathfrak{d} > 0$. On the other hand, for $\delta = 0$ and $p = 2_n^*$, the following nonexistence result was proved in [5].

THEOREM 4.1. If dist(Θ^{Γ} , $\mathbb{R}^{n-1} \times \{0\}$) = dist(Θ , $\mathbb{R}^{n-1} \times \{0\}$), then there exists $\lambda_* \in (-\lambda_1^{\phi}, 0]$ such that, if $\lambda \in (\lambda_*, \infty) \cup \{0\}$, the critical problem $(\wp_{0,2_n}^{\#})$ does not have a least energy solution.

Moreover, $\lambda_* < 0$ if $m \geq 2$.

Proof. See Theorem 3.2 in [5].

For $\delta \in (0, \delta_0)$ and $p \in (2, 2^*_{n-\mathfrak{d}})$, let $J_{\delta,p} : D_0^{1,2}(\Theta_\delta)^{\phi} \to \mathbb{R}$ be the variational functional and $\mathcal{N}^{\phi}_{\delta,p}$ be the Nehari manifold associated to the problem $(\wp_{\delta,p}^{\#})$, and set

$$\ell^{\phi}_{\delta,p} := \inf\{J_{\delta,p}(u) : u \in \mathcal{N}^{\phi}_{\delta,p}\}.$$

We write J_* , \mathcal{N}^{ϕ}_* and ℓ^{ϕ}_* for the variational functional, the Nehari manifold and the infimum associated to the critical problem $(\varphi^{\#}_{0,2^*_n})$ in the whole domain Θ . Extending each function in $\mathcal{N}^{\phi}_{\delta,2^*_n}$ by 0 outside of Θ_{δ} , we have that $\mathcal{N}^{\phi}_{\delta,2^*_n} \subset \mathcal{N}^{\phi}_*$ and $J_{\delta,2^*_n}(u) = J_*(u)$ for every $u \in \mathcal{N}^{\phi}_{\delta,2^*_n}$. Hence, $\ell^{\phi}_* \leq \ell^{\phi}_{\delta,2^*_n}$.

LEMMA 4.2. $\ell^{\phi}_{\delta,2^*} \to \ell^{\phi}_*$ as $\delta \to 0$.

Proof. Let $X := (\mathbb{R}^n)^{\Gamma}$ and Y be its orthogonal complement in \mathbb{R}^n . Since $\Theta \setminus \Theta^{\Gamma} \neq \emptyset$ and every Γ -orbit in $\Theta \setminus \Theta^{\Gamma}$ has positive dimension, we have that $\dim(Y) \geq 2$.

We claim that there are radial functions $\chi_k \in \mathcal{C}_c^{\infty}(Y)$ such that $\chi_k(y) = 1$ if $|y| \leq \frac{1}{k}$,

$$\lim_{k \to \infty} \int_{Y} |\chi_{k}|^{2} = 0 \quad \text{and} \quad \lim_{k \to \infty} \int_{Y} |\nabla \chi_{k}|^{2} = 0.$$
 (12)

To show this, we choose a radial function $g \in C_c^{\infty}(Y)$ such that g(y) = 1 if $|y| \leq 1$ and g(y) = 0 if $|y| \geq 2$, and we set $g_k(y) := g(ky)$. Define

$$\chi_k(y) := \frac{1}{\sigma_k} \sum_{j=1}^k \frac{g_j(y)}{j}, \quad \text{where } \sigma_k := \sum_{j=1}^k \frac{1}{j}.$$

Clearly, $\chi_k(y) = 1$ if $|y| \leq \frac{1}{k}$ and $\chi_k(y) = 0$ if $|y| \geq 2$. As dim $(Y) \geq 2$, we have that $\int_Y |\nabla g_k|^2 \leq \int_Y |\nabla g|^2$. Hence, for some positive constant C,

$$\int_{Y} \left| \nabla \chi_k \right|^2 \le \frac{C}{\sigma_k^2} \sum_{j=1}^k \frac{1}{j^2} \to 0 \qquad \text{as } k \to \infty.$$

Finally, as all functions χ_k are supported in the closed ball of radius 2 in Y, the Poincaré inequality yields

$$\int_{Y} \left| \chi_k \right|^2 \le C \int_{Y} \left| \nabla \chi_k \right|^2 \to 0,$$

and our claim is proved.

Given $\varepsilon > 0$ we choose $\psi \in \mathcal{N}_*^{\phi}$ such that $J_*(\psi) < \ell_*^{\phi} + \frac{\varepsilon}{2}$. For $(x, y) \in X \times Y$, we define $\psi_k(x, y) := (1 - \chi_k(y))\psi(x, y)$. Note that, as χ_k is radial and ψ is is ϕ -equivariant, ψ_k is also ϕ -equivariant. Moreover, the identities (12) easily imply that $\psi_k \to \psi$ in $D_0^{1,2}(\Theta)$. So, for k large enough, there exists $t_k \in (0, \infty)$ such that $\widetilde{\psi}_k := t_k \psi_k \in \mathcal{N}_*^{\phi}$ and $t_k \to 1$. Hence, $\widetilde{\psi}_k \to \psi$ in $D_0^{1,2}(\Theta)$, and we may choose $k_0 \in \mathbb{N}$ such that $J_*(\widetilde{\psi}_{k_0}) < \ell_*^{\phi} + \varepsilon$. Observe that $\operatorname{supp}(\widetilde{\psi}_k) =$ $\operatorname{supp}(\psi_k) \subset \Theta_{\delta}$ if $\delta < \frac{1}{k}$. So $\widetilde{\psi}_k \in \mathcal{N}_{\delta, 2_n^*}^{\phi}$ if $\delta < \frac{1}{k}$. It follows that

$$\ell_*^{\phi} \leq \ell_{\delta,2_n^*}^{\phi} \leq J_{\delta,2_n^*}(\widetilde{\psi}_{k_0}) = J_*(\widetilde{\psi}_{k_0}) < \ell_*^{\phi} + \varepsilon \qquad \forall \delta \in \left(0, \frac{1}{k_0}\right).$$

This finishes the proof.

Set N := n + m and

$$\Omega_{\delta} := \{ (y, z) \in \mathbb{R}^{n-1} \times \mathbb{R}^{m+1} : (y, |z|) \in \Theta_{\delta} \}, \qquad \delta \in [0, \delta_0).$$

Note that Ω_{δ} is $[\Gamma \times O(m+1)]$ -invariant, i.e., $(\gamma y, \varrho z) \in \Omega_{\delta}$ for every $(y, z) \in \Omega_{\delta}$, $\gamma \in \Gamma$, $\varrho \in O(m+1)$. A straightforward computation shows that $u_{\delta,p}$ is a least energy solution to the problem $(\wp_{\delta,p}^{\#})$ if and only if $v_{\delta,p}(y,z) := u_{\delta,p}(y,|z|)$ is a least energy solution to the problem

$$(\wp_{\delta,p}) \qquad \begin{cases} -\Delta v + \lambda v = |v|^{p-2}v & \text{in } \Omega_{\delta}, \\ v = 0 & \text{on } \partial\Omega_{\delta}, \\ v(\gamma y, \varrho z) = \phi(\gamma)v(y, z), \quad \forall \gamma \in \Gamma, \ \varrho \in O(m+1), \ (y, z) \in \Omega_{\delta}. \end{cases}$$

Therefore, for every $\lambda \in (-\lambda_1^{\phi}, \infty)$, $\delta \in (0, \delta_0)$ and $p \in (2, 2^*_{n-\mathfrak{d}})$, the problem $(\wp_{\delta,p})$ has a least energy solution. The following results describe its limit profile.

THEOREM 4.3. For $\delta \in (0, \delta_0)$ let $v_{\delta,*}$ be a least energy solution to the problem $(\wp_{\delta,2^*_{N,m}})$. Assume that

$$\operatorname{dist}(\Theta^{\Gamma}, \mathbb{R}^{n-1} \times \{0\}) = \operatorname{dist}(\Theta, \mathbb{R}^{n-1} \times \{0\}).$$

Then, there exists $\lambda_* \leq 0$ such that, if $\lambda \in (\lambda_*, \infty) \cup \{0\}$, there exist sequences (δ_k) in $(0, \delta_0)$, (ε_k) in $(0, \infty)$, (ζ_k) in Θ^{Γ} , and a nontrivial solution ω to the limit problem (3) such that

(i) $\delta_k \to 0, \ \varepsilon_k^{-1} \operatorname{dist}(\zeta_k, \partial \Theta) \to \infty, \ and \ \zeta_k \to \zeta \ with$

$$\operatorname{dist}(\zeta, \mathbb{R}^{n-1} \times \{0\}) = \operatorname{dist}(\Theta, \mathbb{R}^{n-1} \times \{0\}),$$

- (ii) ω is ϕ -equivariant and has minimal energy among all nontrivial ϕ -equivariant solutions to the problem (3),
- (iii) $v_{\delta_k,*} = \widetilde{\omega}_{\varepsilon_k,\zeta_k} + o(1)$ in $D^{1,2}(\mathbb{R}^N)$, where

$$\widetilde{\omega}_{\varepsilon_k,\zeta_k}(y,z) := \varepsilon_k^{(2-n)/2} \omega\left(\frac{(y,|z|) - \zeta_k}{\varepsilon_k}\right).$$

Moreover, $\lambda_* < 0$ if $m \geq 2$.

Proof. Let λ_* be the number given by Theorem 4.1. Fix $\lambda \in (\lambda_*, \infty) \cup \{0\}$, and let $u_{\delta,*}$ be the least energy solution to the problem $(\wp_{\delta,2_n^*}^{\#})$ given by $v_{\delta,*}(y,z) = u_{\delta,*}(y,|z|)$. Choose a sequence $\delta_k \to 0$ and set $u_k := u_{\delta_k,*}$. Then, $u_k \in \mathcal{N}_*^{\phi}$ and, by Lemma 4.2, $J_*(u_k) \to \ell_*^{\phi}$. It follows from Corollary 3.2 and Theorem 4.1 that, after passing to a subsequence, there exist sequences (ε_k) in $(0,\infty)$ and (ζ_k) in Θ^{Γ} , and a nontrivial ϕ -equivariant solution ω to the limit problem (3) such that $\zeta_k \to \zeta$, $\varepsilon_k^{-1} \operatorname{dist}(\zeta_k, \partial \Theta) \to \infty$, $J_{\infty}(\omega) = \ell_{\infty}^{\phi}$,

$$\lim_{k \to \infty} \left\| u_k - \varepsilon_k^{\frac{2-n}{2}} \omega\left(\frac{\cdot - \zeta_k}{\varepsilon_k}\right) \right\| = 0, \tag{13}$$

and

$$\left[\operatorname{dist}(\zeta, \mathbb{R}^{n-1} \times \{0\})\right] = \min_{x \in \overline{\Theta}} \left[\operatorname{dist}(x, \mathbb{R}^{n-1} \times \{0\})\right].$$

Equation (13) implies that $v_{\delta_k,*}$ satisfies (3). This concludes the proof.

THEOREM 4.4. For $\delta \in (0, \delta_0)$ and $p \in (2^*_{N,m}, 2^*_{N,m+\mathfrak{d}})$ let $v_{\delta,p}$ be a least energy solution to the problem $(\wp_{\delta,p})$. Assume that

$$\operatorname{dist}(\Theta^{\Gamma}, \mathbb{R}^{n-1} \times \{0\}) = \operatorname{dist}(\Theta, \mathbb{R}^{n-1} \times \{0\}).$$

Then, there exists $\lambda_* \leq 0$ such that, if $\lambda \in (\lambda_*, \infty) \cup \{0\}$, there exist sequences (δ_k) in $(0, \delta_0)$, (ε_k) in $(0, \infty)$, (p_k) in $(2^*_{N,m}, 2^*_{N,m+\mathfrak{d}})$, and (ζ_k) in Θ^{Γ} , and a nontrivial solution ω to the limit problem (3) such that

(i)
$$\delta_k \to 0, \ p_k \to 2^*_{N,m}, \ \varepsilon_k^{-1} \text{dist}(\zeta_k, \partial \Theta) \to \infty, \ and \ \zeta_k \to \zeta \ with$$

$$\operatorname{dist}(\zeta, \mathbb{R}^{n-1} \times \{0\}) = \operatorname{dist}(\Theta, \mathbb{R}^{n-1} \times \{0\}),$$

- (ii) ω is ϕ -equivariant and has minimal energy among all nontrivial ϕ -equivariant solutions to the problem (3),
- (iii) $v_{\delta_k,p_k} = \widetilde{\omega}_{\varepsilon_k,\zeta_k} + o(1)$ in $D^{1,2}(\mathbb{R}^N)$, where

$$\widetilde{\omega}_{\varepsilon_k,\zeta_k}(y,z) := \varepsilon_k^{(2-n)/2} \omega\left(\frac{(y,|z|) - \zeta_k}{\varepsilon_k}\right).$$

Moreover, $\lambda_* < 0$ if $m \geq 2$.

Proof. Let λ_* be the number given by Theorem 4.1. Fix $\lambda \in (\lambda_*, \infty) \cup \{0\}$. Let $u_{\delta,p}$ be the least energy solution to the problem $(\wp_{\delta,p}^{\#})$ given by $v_{\delta,p}(y,z) = u_{\delta,p}(y,|z|)$ and let $t_{\delta,p} \in (0,\infty)$ be such that $\widetilde{u}_{\delta,p} := t_{\delta,p}u_{\delta,p} \in \mathcal{N}_{\delta,2_n^*}^{\phi} \subset \mathcal{N}_*^{\phi}$. Proposition 2.5 and Lemma 4.2 allow us to choose $\delta_k \in (0,\delta_0)$ and $p_k \in (2_n^*, 2_{n-\delta}^*)$ such that $\delta_k \to 0$, $p_k \to 2_n^*$, and $J_*(\widetilde{u}_k) \to \ell_*^{\phi}$, where $\widetilde{u}_k := \widetilde{u}_{\delta_k,p_k}$. The rest of the proof is the same as that of Theorem 4.3

Finally, we derive Theorems 1.1 and 1.2 from Theorems 2.3 and 4.4.

Proof of Theorem 1.1. Let $\Gamma := O(n-1)$ and ϕ be the trivial homomorphism $\phi \equiv 1$. Then, $B^{\Gamma} = B \cap [\{0\} \times (0, \infty)]$. A ϕ -equivariant function is simply a Γ -invariant function and, as the standard bubble is radial, it is the least energy Γ -invariant solution to the problem (3), which is unique up to translations and dilations. Since dim $(\Gamma x) = n-2 \geq 1$ for every $x \in B \setminus B^{\Gamma}$, applying Theorems 2.3 and 4.4 to $\Theta := B$ with this group action we obtain Theorem 1.1.

Proof of Theorem 1.2. For $n \geq 5$, let Γ be the subgroup of O(n-1) generated by $\{e^{i\vartheta}, \alpha, \tau : \vartheta \in [0, 2\pi), \alpha \in O(n-5)\}$ acting on a point $y = (\eta, \xi) \in \mathbb{C}^2 \times \mathbb{R}^{n-5} \equiv \mathbb{R}^{n-1}, \eta = (\eta_1, \eta_2) \in \mathbb{C} \times \mathbb{C}$, as

$$e^{i\vartheta}y := (e^{i\vartheta}\eta, \xi), \qquad \alpha y := (\eta, \alpha \xi), \qquad \tau y := (-\overline{\eta}_2, \overline{\eta}_1, \xi),$$

and let ϕ be the homomorphism given by $\phi(e^{i\vartheta}) = 1 = \phi(\alpha)$ and $\phi(\tau) = -1$. Then, $B^{\Gamma} = B \cap [\{0\} \times (0, \infty)]$. If n = 5 then dim $(\Gamma y) = 1$ for every $y \in \mathbb{R}^{n-1} \setminus \{0\}$, whereas for $n \ge 6$ we have that

dim (
$$\Gamma y$$
) =

$$\begin{cases}
n-5 & \text{if } \eta \neq 0 \text{ and } \xi \neq 0, \\
1 & \text{if } \xi = 0, \\
n-6 & \text{if } \eta = 0.
\end{cases}$$

Therefore, if n = 5 or $n \ge 7$, we have that $\dim(\Gamma x) \ge 1$ for every $x \in B \setminus B^{\Gamma}$. Notice that any point $x_0 = (\eta, \xi) \in B$ with $\eta \ne 0$ satisfies condition (5). Hence, Theorem 1.2 follows from Theorems 2.3 and 4.4.

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